

Symmetries & Differential Equations

Applications of Lie Groups

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Motivating question: how can we use symmetries to work with differential equations?

Given a solution, can we find a way to continuously map to other solutions?

We will:

- Describe a structure for describing continuous symmetries
- State a requirement for a differential equation to admit a symmetry

Discrete to Continuous Symmetries



Lie group: **differentiable** manifold with **differentiable** operation, inversion.

↪ Provides a structure for **continuous symmetries**

Ex) Planar rotations $SO(2)$, reals \mathbb{R} under addition

Manifold: looks "pretty close" to a subset of \mathbb{R}^n : circle, sphere, torus, etc

The Lie Algebra

On G live vector fields; functions that assign vectors to points.

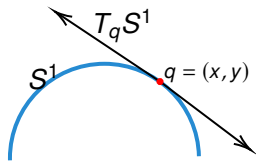
$$\mathbf{v} : M \rightarrow TM$$

$$\mathbf{v}|_x = \xi^1(x)\partial_{x^1} + \cdots + \xi^m(x)\partial_{x^m}$$

Lie algebra of G : $\mathfrak{g} := \{\text{"group operation compatible" vector fields}\}$

Vector fields give rise to **flows**: we denote $\exp(\varepsilon\mathbf{v})x$

$$\mathbf{v}|_x = \left. \frac{d}{d\varepsilon} \exp(\varepsilon\mathbf{v})x \right|_{\varepsilon=0} \quad \exp(\varepsilon\mathbf{v})x = x + \varepsilon\xi(x) + \mathcal{O}(\varepsilon^2)$$



\rightsquigarrow **Linearized**

$$\mathfrak{g} \xrightarrow{\exp} G$$

We now have an **algebraic** means of working with our **group**.
We can replace **non-linear** conditions that arise from the group with
linear conditions in the algebra.

Rather than thinking of how a group operation behaves, we can think
of how the corresponding vector field "flows".

("Because working with a group is cursed" - William)

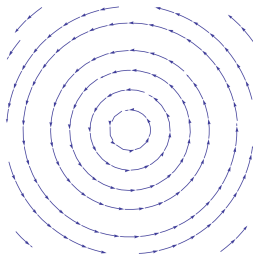
SO(2): Vector Field of a Rotation

$$(x, y) \xrightarrow{g_\theta} (x \cos \theta - y \sin \theta, x \sin \theta + y \cos \theta)$$

$$\mathbf{v}|_{(x,y)} = \left. \frac{d}{d\varepsilon} (x \cos \varepsilon - y \sin \varepsilon, x \sin \varepsilon + y \cos \varepsilon) \right|_{\varepsilon=0} = -y\partial_x + x\partial_y$$

Re-deriving the group operation:

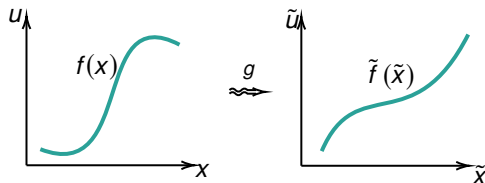
$$\frac{dx}{d\varepsilon} = -y, \quad \frac{dy}{d\varepsilon} = x \implies g_\varepsilon,$$



Extending to Differential Equations

↪ We can find the effect of a group transformation on a function by "flowing" the function on \mathbf{v} ; we work on the space

$(x, u) \in X \times U =$ dependent \times independent



Transforms **functions**, so transforms **derivatives** of functions - how? Formally, we use "prolongation theory" and "prolonged vector fields"

$$\text{pr}^{(n)}\mathbf{v} = \mathbf{v} + \sum_{i=1}^p \phi^{x^i} \partial_{u_{x^i}} + \dots$$

defined over the "jet space" $X \times U^{(n)}$ of $X \times U$.

Infinitesimal Criterion for Invariance

Theorem (Symmetry Condition)

For a differential equation $\Delta(x, u^{(n)}) = 0$ and group of transformations G , if for every infinitesimal generator \mathbf{v} of G ,

$$\text{pr}^{(n)}(\mathbf{v}[\Delta(x, u^{(n)})]) = 0 \quad \text{whenever} \quad \Delta(x, u^{(n)}) = 0,$$

then G is a symmetry group of Δ ; G maps solutions to other solutions.

\rightsquigarrow If I flow a differential equation along a vector field, and the result also satisfies the differential equation, then I have a symmetry of the equation.

"Infinitesimal invariance criterion"

An Aside: Noether's Theorem

Theorem (Noether's First Theorem)

If \mathbf{v} generates a symmetry group of the variational problem $L[u]$, then there exists a corresponding **conservation law** of the Euler-Lagrange equations $\mathcal{E}(L) = 0$.

\rightsquigarrow Connection between **symmetries** and **conservation laws**

If we allow "generalized symmetries"

$$\mathbf{v} = \sum_{i=1}^p \xi_i[u] + \sum_{i=1}^q \phi_i[u]$$

this becomes a **one-to-one** correspondance.

The Heat Equation: Characterization of Symmetry Groups

Given a **differential equation** $\Delta = 0$, how do we find all symmetries of the equation? We consider the heat equation

$$\Delta(t, x, u^{(2)}) = u_{xx} - u_t = 0,$$

and an arbitrary vector field

$$\mathbf{v} = \xi(x, t, u)\partial_x + \tau(x, t, u)\partial_t + \phi(x, t, u)\partial_u,$$

for which the **infinitesimal invariance criterion** (using prolongation) is

$$\phi^t = \phi^{xx}.$$

We expand this to find a system of PDEs for the coefficients of the vector fields \mathbf{v} , and solve to find all independent vector fields.

Heat Equation: Symmetry Groups

Ex)

$$\mathbf{v}_2 = \partial_t \Rightarrow \exp(\varepsilon \mathbf{v})(x, t, u) = (x, t + \varepsilon, u), \text{ Time translation!}$$

$$\mathbf{v}_4 = x\partial_x + 2t\partial_t \Rightarrow \exp(\varepsilon \mathbf{v})(x, t, u) = (e^\varepsilon x, e^{2\varepsilon} t, u), \text{ Space-time scaling!}$$

Symmetries:

G_i	$\exp(\varepsilon \cdot \mathbf{v}_i)(x, t, u)$	Physical Interpretation
1	$(x + \varepsilon, t, u)$	Space Translation
2	$(x, t + \varepsilon, u)$	Time Translation
3	$(x, t, e^\varepsilon u)$	Linearity (Constants)
4	$(e^\varepsilon x, e^{2\varepsilon} t, u)$	Space/Time Scaling
5	$(x + 2\varepsilon t, t, u \cdot e^{-\varepsilon x - \varepsilon^2 t})$	Galilean Boost
6	$(\frac{x}{1 - 4\varepsilon t}, \frac{t}{1 - 4\varepsilon x}, u\sqrt{1 - 4\varepsilon t} e^{\frac{-\varepsilon x^2}{1 - 4\varepsilon t}})$?
α	$(x, t, u + \varepsilon \alpha(x, t))$	Linearity (Additivity)

Given a solution to $\Delta = 0$, applying G_i to it yields another solution.

Heat Equation: Fundamental Solution

Note that $\Delta(c) = 0$ for any constant c .

Applying G_6 with $\varepsilon = 1$ and then G_2 with $\varepsilon = -\frac{1}{4}$ to $u = c := \sqrt{\frac{1}{\pi}}$ we have

$$c \xrightarrow{G_6} \frac{c}{\sqrt{1+4t}} \cdot \exp\left(\frac{-x^2}{1+4t}\right) \xrightarrow{G_2} \frac{1}{\sqrt{4\pi t}} \exp\left\{\frac{-x^2}{4t}\right\},$$

the fundamental solution of the heat equation.

We:

- Formalized continuous symmetries
- Found infinitesimal invariance criterion for symmetry of differential equations
- Applied to the heat equation

↪ **Tools for describing symmetries of differential equations**

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